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Blasius flow and heat transfer of a fourth grade fluid with slip

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Abstract

This investigation deals with the effects of slip, magnetic field and non-Newtonian flow parameters on flow and heat transfer of an incompressible, electrically conducting fourth grade fluid past an infinite porous plate subject to suction and blowing, known as the Blasius flow problem. The heat transfer analysis has been carried out for two heating processes, namely, when the plate is assumed to be at higher temperature than the fluid, and in the second case, the plate is assumed to be insulated. The order of the resulting differential equations exceeds the available adherence boundary conditions due to the presence of the higher order material derivatives of the strain tensor in the constitutive relation for the fourth grade fluid. The system of highly non-linear differential equations is solved using a fourth order Runge-Kutta method along with a shooting method for moderate values of the flow and heat transfer parameters. The effective Broyden's technique has been adopted in order to improve the initial guesses and to satisfy the boundary conditions at infinity. The combined effects of the non-Newtonian flow parameters and slip on the momentum and thermal boundary layers have been discussed in details. An exceptional cross-overs is obtained in the velocity profile in presence of slip. The fourth grade fluid parameter is found to increase the momentum boundary layer thickness, whereas the slip parameter substantially decreases it. Similarly, the non-Newtonian fourth grade fluid parameters and the slip have opposite effects on the thermal boundary layer thickness.

Key words: Blasius flow, Partial slip, Fourth grade fluid, Broyden's method.

1 Introduction

Despite their applications in many industrial processes, the boundary layer flow of higher order fluid of differential type, namely the fourth grade fluid here, has drawn the attention of only few researchers during the last few decades. The main reason is that the constitutive equation of the fourth grade fluid is very complex involving a wide number of parameters. Moreover, the order of the differential equation corresponding to the momentum equation is exceptionally large, and exceeds the number of available boundary conditions, even in the simplest boundary layer flow past a flat plate. Renardy [1] has shown that such fluids are the last to have a stable state of rest. They have been considered in few other configurations such as plane Poiseuille [2,3] or Couette [4] flows or recently, in an asymmetric channel [5].

The unsteady and steady flows of a fourth grade fluid, subject to a transverse uniform magnetic field past a porous plate has been investigated by Hayat *et al.* [6]. The boundary layer approximation resulted in a highly nonlinear differential equation of order six corresponding to the momentum equation with only two adherence boundary conditions. The finite difference method has been adopted to solve the sixth order nonlinear differential equation by augmenting four extra asymptotic boundary conditions. Hayat *et al.* [7] have studied the unsteady flow of a fourth grade fluid past a porous plate. They have successively reduced the order and number of independent variables of the resulting partial differential equations using Lie point symmetries. Subsequently, Sajid *et al.* [8] have studied the steady flow of a fourth grade fluid past a plate with suction and injection velocities at the surface of the plate. The homotopy analysis method (HAM) is used to find the solution of the fifth order highly nonlinear differential equation with two available boundary conditions. Hayat *et al.* [9] have used the HAM method to investigate the influence of the heat transfer on the flow of a fourth grade fluid past a porous plate. The unsteady MHD flow of a fourth grade fluid due to an oscillating plate with suction and blowing has been investigated numerically by Wang and Wu [10]. The corresponding sixth order nonlinear partial differential equation has been solved by finite difference technique after augmenting four asymptotic boundary conditions. Recently, Marinca *et al.* [11] have used optimal homotopy asymptotic method to solve the resulting highly nonlinear momentum equation with no-slip boundary condition.

The foremost reason which motivated for the present study is the nonlinearity and the order of the differential equation corresponding to the momentum equation due to the flow of a fourth grade fluid. The order of the resulting differential equation exceeds drastically from the available adherence boundary conditions due to the presence of the higher order material derivatives of the strain tensor in the constitutive relation for the fourth grade fluid. The

literature survey reveals that in the aforementioned studies, respective authors [6–10] have augmented additional boundary conditions at infinity. However, these are relatively weak and physically plausible. Few other interesting results from the literature survey, which arouse a curiosity to consider the MHD flow and heat transfer of a fourth grade fluid with partial slip boundary condition are as follows. Hayat et al. [6] have reported that in presence of the magnetic field, the boundary layer thickness decreases with an increase in the fourth grade fluid parameter, and the magnetic field has a substantial effect on the boundary layer thickness. The conclusions of the study of the steady flow of a fourth grade fluid past a porous plate by Sajid et al. [8] reveal that there does not exist a solution in the case of injection and the boundary layer thickness is rapidly decreased with an increase in the fourth grade fluid parameter for all values of the suction parameter exceeding one. Moreover, the recent study [9] of the flow and heat transfer of a fourth grade fluid past a porous plate depicts that the thermal boundary layer thickness increases with an increase in the fourth grade fluid parameter. We can cite also the recent contribution of Islam et al. [3], who investigated the steady flow of fourth-grade fluid in a infinite plane Poiseuille channel with slippage using the optimal HAM. They highlighted that this technique provides more satisfactory results than the homotopy perturbation technique.

Hayat et al. [12] considered partial slip effects on the flow and heat transfer characteristics of a third grade fluid past a porous plate using the homotopy analysis method. Nadeem et al. [13] studied the effects of partial slip on a fourth-grade fluid with variable viscosity contained in a stationary cylinder. Analytic solutions of velocity and temperature have been developed and discussed for the Reynolds and Vogels models. The present paper is the extension of the recent work of Sahoo [14], who investigated the flow and heat transfer of an electrically conducting third grade fluid past an infinite plate with partial slip. It deals in the present case with the numerical study of the slip flow and heat transfer of an incompressible, electrically conducting fourth grade fluid past an infinite porous plate subject to suction and blowing. The heat transfer analysis has been carried out for two heating processes, namely, when the plate is assumed to be at higher temperature than the fluid, and in the second case, the plate is assumed to be insulated. To the best of our knowledge, no attention has been given to the effects of partial slip on the MHD boundary layer flow and heat transfer of a fourth grade fluid over a flat plate. Even, the hydrodynamic boundary layer flow and heat transfer of a fourth grade fluid with slip boundary condition is not discussed so far. However, some non-Newtonian fluids like polymer melts exhibit macroscopic wall slip, which makes the no-slip condition fully inappropriate. It is noteworthy to point out that the quantitative behavior of higher grade fluids is strongly dependent on the choice of the material parameters. At the present time, there appears to be no experimental data determining these parameters, particularly the material parameters in the theory of fourth grade fluids. The present investigation is a

step towards understanding the effects of the slip parameter on the flow and heat transfer of a fourth grade fluid.

This paper is divided as follows: in Section 2, the problem is mathematically formulated to get a system of ordinary differential equations with appropriate boundary conditions and is solved using an effective second order numerical scheme in Section 3. The effects of the flow parameters (partial slip, heat transfer, non newtonian parameters, suction/injection velocity) on a Blasius flow of a fourth grade fluid are discussed in details in Section 4 before some concluding remarks in Section 5.

2 Formulation of the problem

We consider the steady two-dimensional flow of an incompressible fourth grade fluid past an infinite porous plate (infinite in the x - direction), coinciding with the plane $y = 0$. The flow far away from the plate is uniform, and is in the direction x parallel to the plate. A transverse uniform magnetic field $(0, B_0, 0)$ (Fig. 1) has been applied at the surface of the plate and the surface admits partial slip.

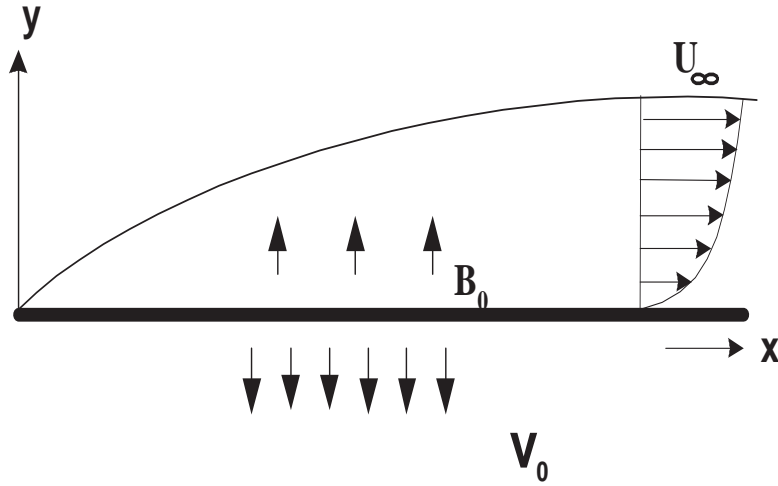


Fig. 1. A sketch of the Blasius flow problem.

Thus, we seek velocity and temperature fields under the form:

$$\mathbf{V} = u(y)\mathbf{i} + v(y)\mathbf{j} \quad (1)$$

$$\theta = \theta(y) \quad (2)$$

where u and v are the components of the velocity in the x and y directions respectively and θ denotes the temperature.

2.1 Flow analysis

Neglecting the body force, the equations of motion and continuity for an electrically conducting, incompressible fluid write:

$$\rho \frac{d\mathbf{V}}{dt} = \nabla \cdot \mathbf{T} + \mathbf{J} \times \mathbf{B} \quad (3)$$

$$\nabla \cdot \mathbf{V} = 0 \quad (4)$$

where d/dt denotes the temporal material derivative, ρ is the density of the fluid assumed to be constant, \mathbf{T} is the Cauchy stress for an incompressible homogeneous fluid of grade four, and \mathbf{J} is the electric current. The term due to the Lorentz force ($\mathbf{J} \times \mathbf{B}$) is simplified by making some assumptions valid for small magnetic Reynolds numbers and no displacement current [15]: the physical quantities are kept constant, the magnetic field is perpendicular to the velocity field, the induced magnetic field remains small compared to the applied magnetic one and the electrical field is assumed to be zero.

Under the consideration of the flow, it follows from (1) and (4) that:

$$v = -V_0 \quad (5)$$

where $V_0 > 0$ is the suction velocity and $V_0 < 0$ is the injection velocity. V_0 is considered as being constant.

The Cauchy stress \mathbf{T} for an incompressible fourth grade fluid has the constitutive equation (see in [5] for example):

$$\begin{aligned} \mathbf{T} = & -p\mathbf{I} + \mu\mathbf{A}_1 + \alpha_1\mathbf{A}_2 + \alpha_2\mathbf{A}_1^2 \\ & + \beta_1\mathbf{A}_3 + \beta_2(\mathbf{A}_2\mathbf{A}_1 + \mathbf{A}_1\mathbf{A}_2) + \beta_3(tr\mathbf{A}_1^2)\mathbf{A}_1 \\ & + \gamma_1\mathbf{A}_4 + \gamma_2(\mathbf{A}_3\mathbf{A}_1 + \mathbf{A}_1\mathbf{A}_3) + \gamma_3\mathbf{A}_2^2 + \gamma_4(\mathbf{A}_2\mathbf{A}_1^2 + \mathbf{A}_1^2\mathbf{A}_2) \\ & + \gamma_5(tr\mathbf{A}_2)\mathbf{A}_2 + \gamma_6(tr\mathbf{A}_2)\mathbf{A}_1^2 + [\gamma_7(tr\mathbf{A}_3) + \gamma_8 tr(\mathbf{A}_2\mathbf{A}_1)]\mathbf{A}_1 \end{aligned} \quad (6)$$

where \mathbf{I} is the identity tensor. The coefficient μ is the coefficient of viscosity ($kg/m/s$). α_1 and α_2 are the second order material constants (kg/m), β_1, \dots, β_3 the third order material constants ($kg.s/m$) and $\gamma_1, \dots, \gamma_8$ the fourth order material constants ($kg.s^2/m$). The first to fourth order Rivlin-Ericksen tensors \mathbf{A}_i ($i = 1, \dots, 4$) are defined by the recursive relation:

$$\mathbf{A}_0 = \mathbf{I} \quad (7)$$

$$\mathbf{A}_n = \frac{d\mathbf{A}_{n-1}}{dt} + \mathbf{A}_{n-1} \nabla \mathbf{V} + (\nabla \mathbf{V})^T \mathbf{A}_{n-1}, n \geq 1 \quad (8)$$

where $(\nabla \mathbf{V})^T$ is the transpose of $\nabla \mathbf{V}$.

Now substituting (1) and (5) in the constitutive equation (6) for the fourth grade fluid, we obtain the following physical components of the Cauchy stress tensor \mathbf{T} :

$$T_{xx} = -p + \alpha_2 \left(\frac{du}{dy} \right)^2 - 2\beta_2 V_0 \frac{du}{dy} \frac{d^2 u}{dy^2} + 2\gamma_2 V_0^2 \frac{du}{dy} \frac{d^3 u}{dy^3} + \gamma_3 V_0^2 \left(\frac{d^2 u}{dy^2} \right)^2 + 2\gamma_6 \left(\frac{du}{dy} \right)^4 \quad (9)$$

$$\begin{aligned} T_{xy} = & \mu \frac{du}{dy} - \alpha_1 V_0 \frac{d^2 u}{dy^2} + \beta_1 V_0^2 \frac{d^3 u}{dy^3} + 2(\beta_1 + \beta_3) \left(\frac{du}{dy} \right)^3 - \gamma_1 V_0^3 \frac{d^4 u}{dy^4} \\ & - (6\gamma_2 + 2\gamma_3 + 2\gamma_4 + 2\gamma_5 + 6\gamma_7 + 2\gamma_8) V_0 \left(\frac{du}{dy} \right)^2 \frac{d^2 u}{dy^2} \end{aligned} \quad (10)$$

$$\begin{aligned} T_{yy} = & -p + (2\alpha_1 + \alpha_2) \left(\frac{du}{dy} \right)^2 - 2V_0(3\beta_1 + \beta_2) \frac{du}{dy} \frac{d^2 u}{dy^2} + 6\gamma_1 V_0^2 \frac{d}{dy} \left(\frac{du}{dy} \frac{d^2 u}{dy^2} \right) \\ & + 2(\gamma_1 + \gamma_2) V_0^2 \frac{d^3 u}{dy^3} \frac{du}{dy} + \gamma_3 V_0^2 \left(\frac{d^2 u}{dy^2} \right)^2 + 2(2\gamma_3 + 2\gamma_4 + 2\gamma_5 + \gamma_6) \left(\frac{du}{dy} \right)^4 \end{aligned} \quad (11)$$

$$T_{zz} = -p \quad (12)$$

$$T_{xz} = T_{zy} = 0 \quad (13)$$

$$T_{xy} = T_{yx}, \quad T_{xz} = T_{zx}, \quad T_{yz} = T_{zy} \quad (14)$$

Now, the three scalar components of the momentum equation (3), in view of Eqn. (5), become:

$$-\rho V_0 \frac{du}{dy} = \frac{\partial T_{xx}}{\partial x} + \frac{\partial T_{xy}}{\partial y} + \frac{\partial T_{xz}}{\partial z} - \sigma B_0^2 u \quad (15)$$

$$0 = \frac{\partial T_{xy}}{\partial x} + \frac{\partial T_{yy}}{\partial y} + \frac{\partial T_{yz}}{\partial z} + \sigma B_0^2 V_0 \quad (16)$$

$$0 = \frac{\partial T_{xz}}{\partial x} + \frac{\partial T_{yz}}{\partial y} + \frac{\partial T_{zz}}{\partial z} \quad (17)$$

where σ is the electrical conductivity.

Inserting the above stress components in Eqns. (15)- (17), we obtain:

$$\begin{aligned} \mu \frac{d^2 u}{dy^2} + \rho V_0 \frac{du}{dy} - \alpha_1 V_0 \frac{d^3 u}{dy^3} + \beta_1 V_0^2 \frac{d^4 u}{dy^4} + 6(\beta_2 + \beta_3) \left(\frac{du}{dy} \right)^2 \frac{d^2 u}{dy^2} - \gamma_1 V_0^3 \frac{d^5 u}{dy^5} \\ - 2(3\gamma_2 + \gamma_3 + \gamma_4 + \gamma_5 + 3\gamma_7 + \gamma_8) V_0 \frac{d}{dy} \left[\left(\frac{du}{dy} \right)^2 \frac{d^2 u}{dy^2} \right] - \sigma B_0^2 u = \frac{\partial \hat{p}}{\partial x} \end{aligned} \quad (18)$$

$$\sigma B_0^2 V_0 = \frac{\partial \hat{p}}{\partial y} \quad (19)$$

$$0 = \frac{\partial \hat{p}}{\partial z} \quad (20)$$

where the modified pressure \hat{p} is given by:

$$\begin{aligned} \hat{p} = -T_{yy} = p - (2\alpha_1 + \alpha_2) \left(\frac{du}{dy} \right)^2 + 2V_0(3\beta_1 + \beta_2) \frac{du}{dy} \frac{d^2 u}{dy^2} - 6\gamma_1 V_0^2 \frac{d}{dy} \left(\frac{du}{dy} \frac{d^2 u}{dy^2} \right) \\ - 2(\gamma_1 + \gamma_2) V_0^2 \frac{d^3 u}{dy^3} \frac{du}{dy} - \gamma_3 V_0^2 \left(\frac{d^2 u}{dy^2} \right)^2 - 2(2\gamma_3 + 2\gamma_4 + 2\gamma_5 + \gamma_6) \left(\frac{du}{dy} \right)^4 \end{aligned} \quad (21)$$

Eqn. (20) indicates that \hat{p} is not a function of z . On eliminating the pressure gradient from (18) and (19) by cross differentiation, we finally obtain:

$$\begin{aligned} -\rho V_0 \frac{d^2 u}{dy^2} = \mu \frac{d^3 u}{dy^3} - \sigma B_0^2 \frac{du}{dy} - \alpha_1 V_0 \frac{d^4 u}{dy^4} + \beta_1 V_0^2 \frac{d^5 u}{dy^5} + 6(\beta_2 + \beta_3) \frac{d}{dy} \left[\left(\frac{du}{dy} \right)^2 \frac{d^2 u}{dy^2} \right] \\ - \gamma_1 V_0^3 \frac{d^6 u}{dy^6} - 2(3\gamma_2 + \gamma_3 + \gamma_4 + \gamma_5 + 3\gamma_7 + \gamma_8) V_0 \frac{d^2}{dy^2} \left[\left(\frac{du}{dy} \right)^2 \frac{d^2 u}{dy^2} \right] \end{aligned} \quad (22)$$

Equation (22) is the stationary form of the one used by Wang and Wu [10]. The appropriate partial slip boundary conditions are:

$$\begin{aligned} u(0) &= \lambda_1 T_{xy} \Big|_{y=0} \\ u(y) &\rightarrow U_\infty \quad \text{as } y \rightarrow \infty \end{aligned} \quad (23)$$

where λ_1 is the slip factor and U_∞ is the mainstream velocity.

By introducing the following non-dimensional quantities, with ν the kinematic viscosity of the fluid,

$$\begin{aligned} \bar{u} &= \frac{u}{U_\infty}, \quad \bar{y} = \frac{U_\infty y}{\nu}, \quad \bar{V}_0 = \frac{V_0}{U_\infty}, \\ \bar{M}_n &= \frac{\nu \sigma B_0^2}{\rho U_\infty^2}, \quad \bar{\alpha}_1 = \frac{\alpha_1 U_\infty^2}{\rho \nu^2}, \quad \bar{\beta}_1 = \frac{\beta_1 U_\infty^4}{\rho \nu^3}, \quad \bar{\gamma}_1 = \frac{\gamma_1 U_\infty^6}{\rho \nu^4}, \\ \bar{\beta} &= \frac{6(\beta_2 + \beta_3) U_\infty^4}{\rho \nu^3}, \quad \bar{\gamma} = 2(3\gamma_2 + \gamma_3 + \gamma_4 + \gamma_5 + 3\gamma_7 + \gamma_8) \frac{U_\infty^6}{\rho \nu^4} \end{aligned} \quad (24)$$

the boundary value problem (22) takes the form:

$$\begin{aligned} M_n \frac{du}{dy} - V_0 \frac{d^2 u}{dy^2} - \frac{d^3 u}{dy^3} + \alpha_1 V_0 \frac{d^4 u}{dy^4} - \beta_1 V_0^2 \frac{d^5 u}{dy^5} + \gamma_1 V_0^3 \frac{d^6 u}{dy^6} &= \beta \left[2 \frac{du}{dy} \left(\frac{d^2 u}{dy^2} \right)^2 \right. \\ &\quad \left. + \left(\frac{du}{dy} \right)^2 \frac{d^3 u}{dy^3} \right] - \gamma V_0 \left[2 \left(\frac{d^2 u}{dy^2} \right)^3 + 6 \frac{du}{dy} \frac{d^2 u}{dy^2} \frac{d^3 u}{dy^3} + \left(\frac{du}{dy} \right)^2 \frac{d^4 u}{dy^4} \right] \end{aligned} \quad (25)$$

and the boundary conditions (23) become:

$$\begin{aligned} u &= \lambda \left[\frac{du}{dy} - V_0 \alpha_1 \frac{d^2 u}{dy^2} + V_0^2 \beta_1 \frac{d^3 u}{dy^3} + \frac{\beta}{3} \left(\frac{du}{dy} \right)^3 - V_0^3 \gamma_1 \frac{d^4 u}{dy^4} - V_0 \gamma \left(\frac{du}{dy} \right)^2 \frac{d^2 u}{dy^2} \right] \text{ at } y = 0 \\ u &\rightarrow 1 \quad \text{as } y \rightarrow \infty \end{aligned} \quad (26)$$

where λ is the non-dimensional slip factor indicating the relative importance of the partial slip, and for simplicity, the bars of the non-dimensional quantities have been omitted. M_n is the magnetic parameter and $\alpha_1, \beta_1, \gamma_1, \beta$ and γ are the non-Newtonian flow parameters.

2.2 Heat transfer analysis

We now turn our attention to the energy equation subject to two different sets of boundary conditions. In the first instance, we assume that the plate is at a temperature higher than the fluid. For the second case, we assume that the plate is insulated. Neglecting the radiant heating, the energy equation can be written as:

$$\rho c_p \frac{d\theta}{dt} = \kappa \nabla^2 \theta + \mathbf{T} \cdot \nabla \mathbf{V} \quad (27)$$

where c_p is the mass specific heat ($J/K/kg$) and κ the thermal conductivity ($kg.m/s^3/K$). Since the velocity and temperature depend only on y , using Eqns. (1) and (2) in the above energy equation (27), we get:

$$-\rho c_p V_0 \frac{d\theta}{dy} = \kappa \frac{d^2 \theta}{dy^2} + T_{xy} \frac{du}{dy} \quad (28)$$

The boundary conditions depend on the following two heating processes.

2.2.1 Constant surface temperature

In this case, the appropriate boundary conditions are:

$$\begin{aligned} \theta(0) &= \theta_w, \\ \theta(y) &\rightarrow \theta_\infty \quad \text{as } y \rightarrow \infty \end{aligned} \quad (29)$$

where θ_w is the wall temperature and θ_∞ is the temperature of the ambient fluid.

Using the non-dimensional temperature $\bar{\theta} = \frac{\theta - \theta_\infty}{\theta_w - \theta_\infty}$, the non-dimensional quantities (24) and the expression for T_{xy} from Eqn. (10) in the energy equation (28), we obtain the following non-dimensional highly non-linear differential equation:

$$\begin{aligned} \frac{d^2\theta}{dy^2} + P_r V_0 \frac{d\theta}{dy} + P_r E_c \left[\left(\frac{du}{dy} \right)^2 - \alpha_1 V_0 \frac{d^2u}{dy^2} \frac{du}{dy} + \beta_1 V_0^2 \frac{d^3u}{dy^3} \frac{du}{dy} + \beta \left(\frac{du}{dy} \right)^4 \right. \\ \left. - \gamma V_0 \left(\frac{du}{dy} \right)^3 \frac{d^2u}{dy^2} - \gamma_1 V_0^3 \frac{d^4u}{dy^4} \frac{du}{dy} \right] = 0 \end{aligned} \quad (30)$$

where $P_r = \frac{\mu c_p}{\kappa}$ is the Prandtl number and $E_c = \frac{U_\infty^2}{c_p(T_w - T_\infty)}$ is the Eckert number and bars have been suppressed.

The boundary conditions for θ are:

$$\begin{aligned} \theta(0) &= 1, \\ \theta &\rightarrow 0 \quad \text{as } y \rightarrow \infty \end{aligned} \quad (31)$$

2.2.2 Insulated surface

In this case, the relevant boundary conditions are:

$$\begin{aligned} \frac{d\theta}{dy} \Big|_{y=0} &= 0, \\ \theta(y) &\rightarrow \theta_\infty \quad \text{as } y \rightarrow \infty \end{aligned} \quad (32)$$

Again, using the non-dimensional temperature $\bar{g} = \frac{\theta - \theta_\infty}{\theta_w - \theta_\infty}$, the non-dimensional quantities (24) and the expression for T_{xy} from Eqn. (10) in the energy equation (28), we obtain the following highly non-linear differential equation:

$$\begin{aligned} \frac{d^2g}{dy^2} + P_r V_0 \frac{dg}{dy} + P_r E_c \left[\left(\frac{du}{dy} \right)^2 - \alpha_1 V_0 \frac{d^2u}{dy^2} \frac{du}{dy} + \beta_1 V_0^2 \frac{d^3u}{dy^3} \frac{du}{dy} + \beta \left(\frac{du}{dy} \right)^4 \right. \\ \left. - \gamma V_0 \left(\frac{du}{dy} \right)^3 \frac{d^2u}{dy^2} - \gamma_1 V_0^3 \frac{d^4u}{dy^4} \frac{du}{dy} \right] = 0 \end{aligned} \quad (33)$$

which is same as Eqn.(30) and all the other parameters are the same as before. Bars have been omitted.

The boundary conditions for g are:

$$\begin{aligned} \frac{dg}{dy}\bigg|_{y=0} &= 0, \\ g &\rightarrow 0 \quad \text{as } y \rightarrow \infty \end{aligned} \tag{34}$$

3 Numerical solution of the problem

The system of highly non-linear differential equations (25) and (30) subject to boundary conditions (26) and (31), and the system of equations (25) and (33) subject to boundary conditions (26) and (34) are solved using a shooting method along with the fourth order Runge-Kutta method for different moderate values of the flow and heat transfer parameters. The effective Broyden's technique has been adopted in order to improve the initial guesses and to satisfy the boundary conditions at infinity. Moreover, it eases the computations by avoiding to compute the Jacobian matrix in each integration step. In fact, the above system of differential equations are highly non-linear, whose order is much more higher than the available known boundary conditions, which debars the use of the finite difference method without augmenting asymptotic boundary conditions at infinity [6,10]. It is noteworthy to mention that the shooting technique in calculating process does not require the entire starting boundary conditions, instead it only needs to match with the other end of the boundary condition to add amendment to its starting boundary conditions. Thus, it can process higher order differential equation problem that lacks of sufficient boundary conditions. Despite of few shortcomings of the shooting method, we found the shooting method quite suitable and aesthetically pleasing for the present investigation.

The semi-infinite integration domain $y \in [0, \infty)$ is replaced by a finite domain $y \in [0, y_\infty)$. In fact, y_∞ usually depends upon the physical parameters of the problem and its value needs to be adjusted as the values of the parameters change. In practice, y_∞ should be chosen sufficiently large so that the numerical solution closely approximates the terminal boundary conditions at ∞ . Here, y_∞ has been fixed to 10.

The above systems of differential equations are converted into an equivalent initial value problem in order to use the fourth order Runge-Kutta method. For the sake of brevity, here we have explained the solution procedure for the system (25) and (30) subject to boundary conditions (26) and (31) concisely. The solution procedure for the system (25) and (33) subject to boundary conditions (26) and (34) is just similar.

To reduce the equations to equivalent first order system, we set:

$$\begin{aligned} u &= z_1, & \frac{du}{dy} &= z_2, & \frac{d^2u}{dy^2} &= z_3, & \frac{d^3u}{dy^3} &= z_4, & \frac{d^4u}{dy^4} &= z_5, & \frac{d^5u}{dy^5} &= z_6, \\ \theta &= z_7, & \frac{d\theta}{dy} &= z_8, \end{aligned} \quad (35)$$

to finally get:

$$\begin{aligned} z'_1 &= z_2, & z_1(0) &= \lambda[s_1 - \alpha_1 V_0 s_2 + \beta_1 V_0^2 s_3 + \frac{\beta}{3} s_1^3 - \gamma_1 V_0^3 s_4 - \gamma V_0 s_1^2 s_2] \\ z'_2 &= z_3, & z_2(0) &= s_1 \\ z'_3 &= z_4, & z_3(0) &= s_2 \\ z'_4 &= z_5, & z_4(0) &= s_3 \\ z'_5 &= z_6, & z_5(0) &= s_4 \\ z'_6 &= -M_n z_2 + V_0 z_3 + z_4 - \alpha_1 V_0 z_5 + \beta_1 V_0^2 z_6 + \beta[2z_2 z_3^2 + z_2^2 z_4] \\ &\quad - \gamma V_0[2z_3^3 + 6z_2 z_3 z_4 + z_2^2 z_5], & z_6(0) &= s_5 \\ z'_7 &= z_8, & z_7(0) &= 1 \\ z'_8 &= -P_r V_0 z_8 - P_r E_c[z_2^2 - \alpha_1 V_0 z_3 z_2 + \beta_1 V_0^2 z_4 z_2 + \beta z_2^4 - \gamma V_0 z_2^3 z_3 - \gamma_1 V_0^3 z_2 z_5], \\ &\quad z_8(0) &= s_6. \end{aligned} \quad (36)$$

Here, the prime indicates differentiation with respect to y . The essence of this method is to reduce the boundary value problem to an equivalent initial value problem and then use a shooting numerical technique to guess s_i , $i = 1, 2, \dots, 6$, until the terminal boundary conditions $z_1 = 1$, $z_7 = 0$ are satisfied at $y = y_\infty$. Because of the non-linearities, which appear in the differential equations, the solution by the above method has been iterated and with reasonably close initial guesses to start the iterations, the convergence to the actual values within an accuracy of $O(10^{-6})$ is attained after only 9 to 11 iterations.

4 Results and discussion

Keeping in view of a proliferation of the number of parameters, we have restricted the discussions of the obtained results to the effects of emerging important flow parameters α_1 , β , γ , V_0 , M_n , P_r , E_c , and the newly introduced slip parameter λ on the velocity and temperature fields. It is found that β , β_1 , and similarly γ , γ_1 have similar effects on the velocity and temperature fields. Further, one can refer the works of Hayat et al. [6,9] and Sajid et al. [8] for a precise discussion with the conventional no-slip boundary conditions.

As a preliminary test [16] to check the accuracy of the solution, it may be noted that for the no-slip and non magnetic case, at $y = 0$, Eqn. (25) reduces to:

$$V_0 \frac{d^2 u}{dy^2} - \frac{d^3 u}{dy^3} + \alpha_1 V_0 \frac{d^4 u}{dy^4} - \beta_1 V_0^2 \frac{d^5 u}{dy^5} + \gamma_1 V_0^3 \frac{d^6 u}{dy^6} - \beta \left[2 \frac{du}{dy} \left(\frac{d^2 u}{dy^2} \right)^2 + \left(\frac{du}{dy} \right)^2 \frac{d^3 u}{dy^3} \right] + \gamma V_0 \left[2 \left(\frac{d^2 u}{dy^2} \right)^3 + 6 \frac{du}{dy} \frac{d^2 u}{dy^2} \frac{d^3 u}{dy^3} + \left(\frac{du}{dy} \right)^2 \frac{d^4 u}{dy^4} \right] \Big|_{y=0} = 0. \quad (37)$$

For all solutions computed for different ranges of the flow parameters, the left hand side of Eqn. (37) is found to be less than 10^{-4} .

The variations of the non-dimensional velocity u with different flow parameters have been plotted in Figs. 2-7 for $\beta_1 = 1$ and $\gamma_1 = 1$. Fig. 2 depicts that in presence of slip, as the normal stress coefficient α_1 increases from zero, the velocity increases near the surface and then starts decreasing away from the surface of the plate, resulting in a ‘cross-over’ in the profile. It is noteworthy to mention that such a cross-over is absent for the no-slip case [9]. Figs. 3 and 4 show that velocity decreases, or in other words, the momentum boundary layer thickness increases with an increase in the non-Newtonian fluid parameters β and γ respectively. These findings are in agreement with those reported by Hayat et al. [9] for the flow past a porous flat plate or by Islam et al. [3] for a plane Poiseuille flow. In general, the fourth grade non-Newtonian fluid exhibits a thicker boundary layer than the Newtonian fluid, and increasing the higher order material parameters of non-Newtonian fluid causes further thickening of the boundary layer. Fig. 5 shows the variation of u with the non-dimensional suction velocity. It is again interesting to find that in presence of slip, as the suction increases, the velocity near the plate increases, and then decreases gradually. This results in a cross-over in the velocity profile. However, in absence of slip, such a cross-over is not found, as reported by Sajid et al. [8] and Hayat et al. [9]. Fig. 6 reveals that an increase in the magnetic interaction parameter M_n increases the velocity throughout the domain of integration. The effect of slip on the velocity profile, when other flow parameters are kept constant, can be seen from Fig. 7. As was expected, an increase in slip results in an increase in the velocity near the surface of the plate. The velocity approaches towards the free stream value 1 with an increase in the slip, and flow behaves as though it were inviscid.

Another interesting quantity is the skin-friction coefficient at the wall. It is related to the wall shear stress, given by the Eqn. (10), which in the dimensionless form becomes:

$$T_{xy}|_{y=0} = \left[\frac{du}{dy} - V_0 \alpha_1 \frac{d^2 u}{dy^2} + V_0^2 \beta_1 \frac{d^3 u}{dy^3} + \frac{\beta}{3} \left(\frac{du}{dy} \right)^3 - V_0^3 \gamma_1 \frac{d^4 u}{dy^4} - V_0 \gamma \left(\frac{du}{dy} \right)^2 \frac{d^2 u}{dy^2} \right]_{y=0} \quad (38)$$

The effect of the fourth grade fluid parameter γ on the non-dimensional wall shear stress $T_{xy}|_{y=0}$ is clear from Fig 8. It shows that $T_{xy}|_{y=0}$ increases, resulting in an increase in the momentum boundary layer thickness, with an increase in γ . On the other hand, as was expected, $T_{xy}|_{y=0}$ decreases rapidly with an increase in the slip parameter λ (Fig 9), indicating a substantial decrease in the momentum boundary layer thickness.

Figs. 10-21 indicate the effects of different emerging flow parameters on the non-dimensional temperature profiles $\theta(y)$ and $g(y)$ for the constant surface temperature and the insulated wall cases, when $\beta_1 = \gamma_1 = 1$, $P_r = 3$, and $E_c = 0.5$. Figs 10-12 respectively depict that even in presence of slip, an increase in the non-Newtonian parameters α_1 , β and γ increase the temperature $\theta(y)$. These findings are in agreement with those reported by Hayat et al. [9]. The effect of the fourth grade fluid parameter γ on θ is not prominent throughout the domain of integration, as shown in Fig. 12. Figs. 13-15 show that the temperature, and hence the thermal boundary layer thickness, decreases with an increase in the suction, magnetic parameter and slip parameter respectively. The suction has a prominent effect on the temperature profile. Figs. 16-21 depict the effects of various flow parameters on the non-dimensional temperature $g(y)$ for the insulated wall case. It is clear that all the flow parameters have similar effects of $\theta(y)$ and $g(y)$. Fig. 21 shows that the slip has a substantial effect on the thermal boundary layer for the insulated wall case.

The dimensionless surface temperature gradient $\theta'(0)$ in the constant surface temperature case, and the dimensionless surface temperature $g(0)$ in the insulated wall case against different flow parameters have been tabulated in Table 1. It is clear that the shear thickening parameter β decreases the magnitude of the rate of heat transfer $\theta'(0)$, when the other parameters are kept constant. It can be of interest to find that the heat transfer rate increases with γ in presence of comparatively large value ($\lambda = 2$) of the slip parameter. The slip parameter λ increase the rate of heat transfer at the surface of the plate. Moreover, it is observed that for constant E_c , the value of $|\theta'(0)|$ increases with P_r , since the thermal boundary layer thickness δ_T decreases with P_r due to the estimate $\delta_T \sim \frac{1}{\sqrt{P_r}}$. Apparently, the surface temperature $g(0)$ in the insulated wall case increases with E_c , when other flow parameters are kept constant.

5 Conclusions

Flow and heat transfer of an electrically conducting fourth grade fluid past an infinite plate with partial slip boundary condition, subject to uniform suction has been considered in this paper. Due to the presence of higher order material derivatives of the strain tensor in the constitutive equation of the fourth grade fluid, the momentum equation gives rise to a highly nonlinear differential equation of order six with only two adherence boundary conditions. The resulting system of differential equations with inadequate boundary conditions has been solved by the shooting method along with the fourth order Runge-Kutta method and the Broyden's method. The use of the Broyden's method as a zero-finding algorithm has enhanced the efficiency of the shooting technique by reducing the computational time. The combined effects of the emerging non-Newtonian flow parameters and slip on the momentum and thermal boundary layers have been discussed elaborately and shown graphically. It is interesting to find the exceptional cross-overs in the velocity profile in the presence of slip. The fourth grade fluid parameter γ increases the momentum boundary layer thickness, whereas the slip parameter substantially decreases it. Similarly, the non-Newtonian fourth grade fluid parameters and the slip have opposite effects on the thermal boundary layer thickness.

The present study will be extended in a close future to take into account additional effects such as the dependance of the slip velocity on the normal stress, Joule heating in the energy equation or the thermal slip condition. As the shooting method limited the present investigation to moderate values of the flow parameters, a more effective and suitable numerical technique would be also suitable to see the effects of higher values of the fourth grade and slip parameters on the Blasius flow problem.

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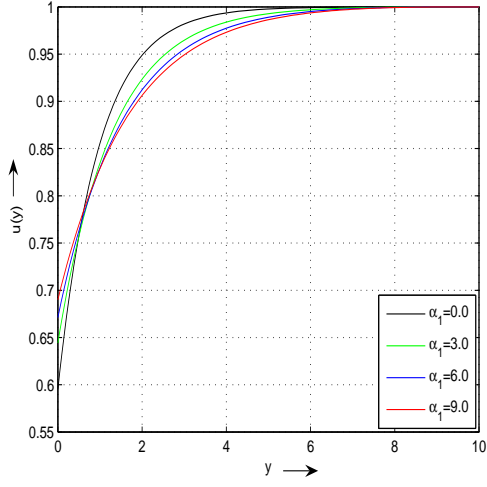


Fig. 2. Variation of u with α_1 at $\beta = 2.0$, $\gamma = 1.0$, $V_0 = 0.5$, $M_n = 1.0$ & $\lambda = 1.0$.

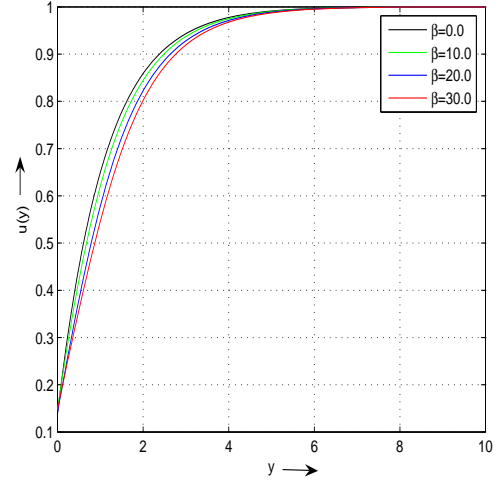


Fig. 3. Variation of u with β at $\alpha_1 = 1.0$, $\gamma = 1.0$, $V_0 = 0.5$, $M_n = 1.0$ & $\lambda = 0.1$.

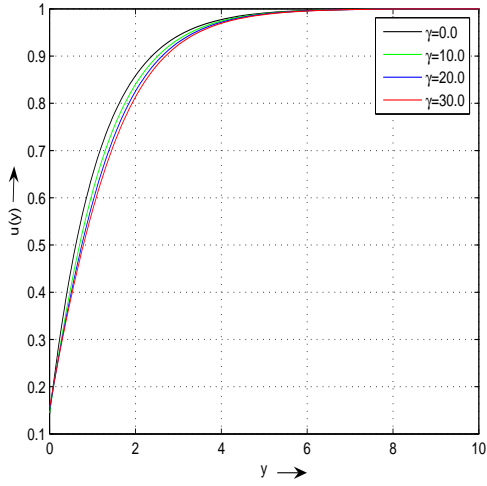


Fig. 4. Variation of u with γ at $\alpha_1 = 1.0$, $\beta = 2.0$, $V_0 = 0.5$, $M_n = 1.0$ & $\lambda = 0.1$.

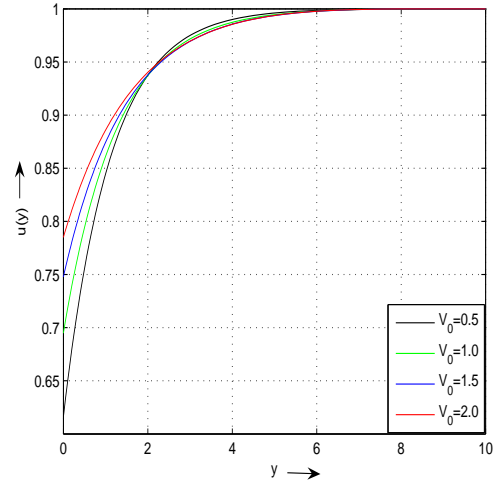


Fig. 5. Variation of u with V_0 at $\alpha_1 = 1.0$, $\beta = 2.0$, $\gamma = 1.0$, $M_n = 1.0$ & $\lambda = 1.0$.

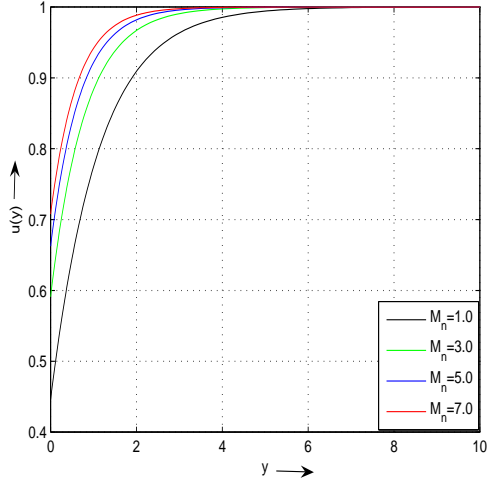


Fig. 6. Variation of u with M_n at $\alpha_1 = 1$, $\beta = 2$, $\gamma = 1$, $V_0 = 0.5$ & $\lambda = 0.5$.

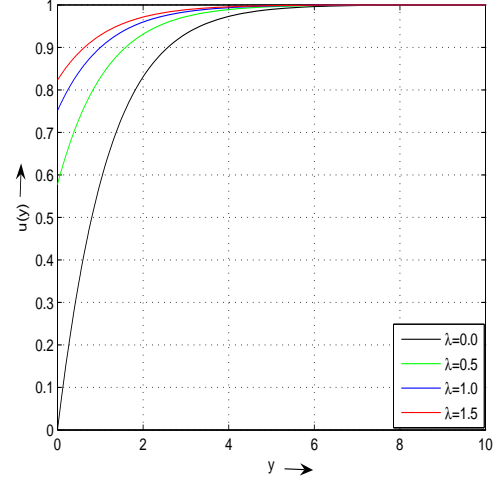


Fig. 7. Variation of u with λ at $\alpha_1 = 1$, $\beta = 2$, $\gamma = 1$, $V_0 = 0.5$ & $M_n = 1$.

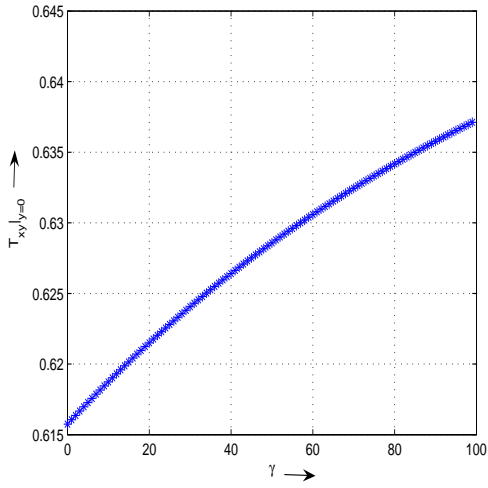


Fig. 8. Variation of $T_{xy}|_{y=0}$ with γ at $\alpha_1 = 1$, $\beta = 2$, $V_0 = 0.5$, $M_n = 1$ & $\lambda = 1$.

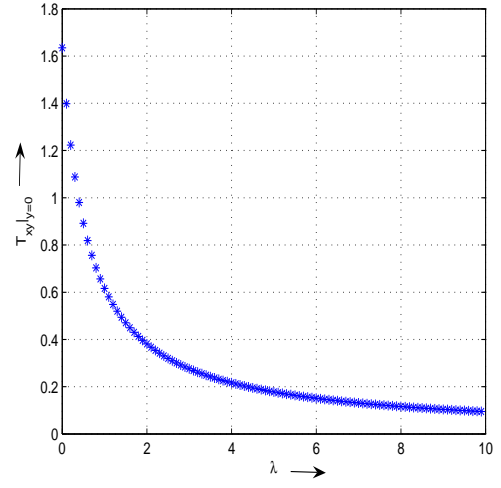


Fig. 9. Variation of $T_{xy}|_{y=0}$ with λ at $\alpha_1 = 1$, $\beta = 2$, $\gamma = 1$, $V_0 = 0.5$ & $M_n = 1$.

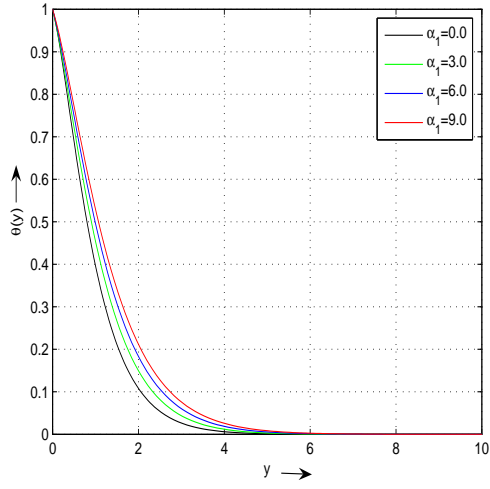


Fig. 10. Variation of θ with α_1 at $\beta = 2$, $\gamma = 1$, $V_0 = 0.5$, $M_n = 1$ & $\lambda = 0.1$.

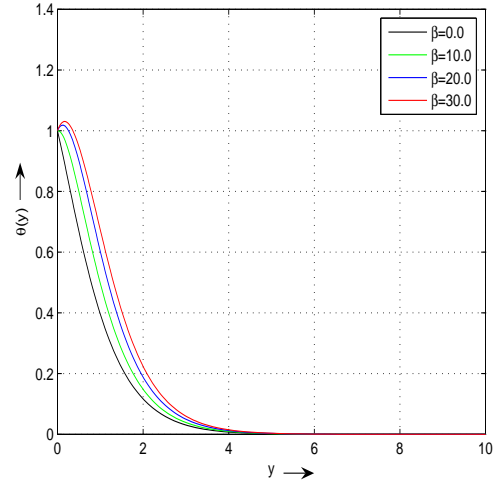


Fig. 11. Variation of θ with β at $\alpha_1 = 1$, $\gamma = 1$, $V_0 = 0.5$, $M_n = 1$ & $\lambda = 0.1$.

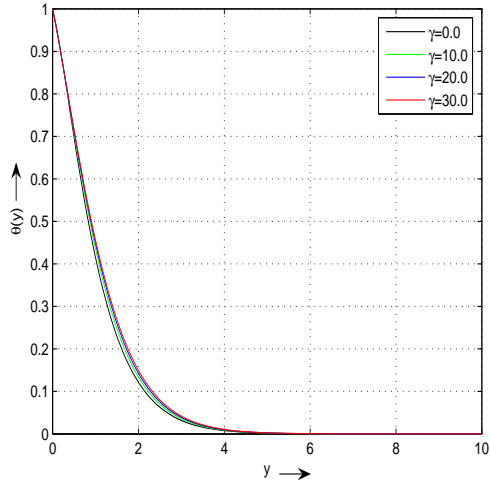


Fig. 12. Variation of θ with γ at $\alpha_1 = 1$, $\beta = 2$, $V_0 = 0.5$, $M_n = 1$ & $\lambda = 0.1$.

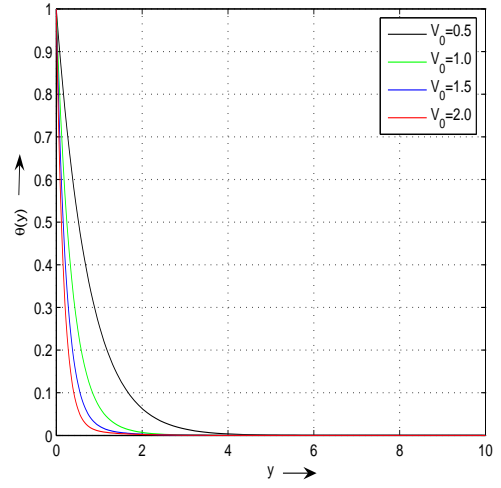


Fig. 13. Variation of θ with V_0 at $\alpha_1 = 1$, $\beta = 2$, $\gamma = 1$, $M_n = 1$ & $\lambda = 1$.

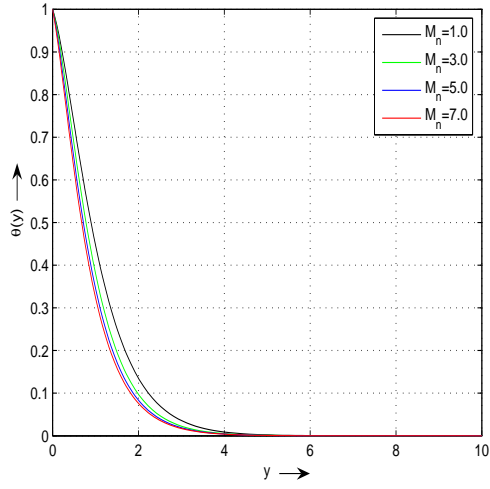


Fig. 14. Variation of θ with M_n at $\alpha_1 = 1$, $\beta = 2$, $\gamma = 1$, $V_0 = 0.5$ & $\lambda = 0.5$.

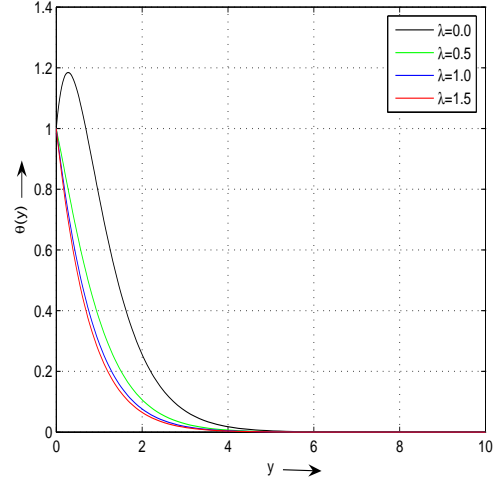


Fig. 15. Variation of θ with λ at $\alpha_1 = 1$, $\beta = 2$, $\gamma = 1$, $V_0 = 0.5$ & $M_n = 1$.

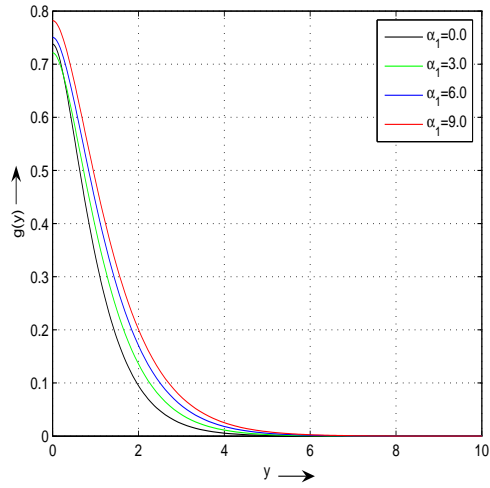


Fig. 16. Variation of g with α_1 at $\beta = 2$, $\gamma = 1$, $V_0 = 0.5$, $M_n = 1$ & $\lambda = 0.1$.

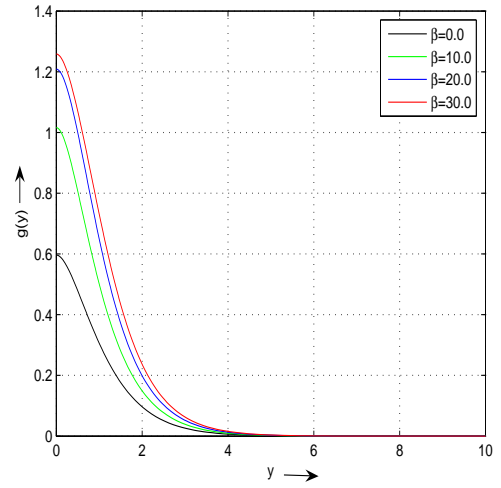


Fig. 17. Variation of g with β at $\alpha_1 = 1$, $\gamma = 1$, $V_0 = 0.5$, $M_n = 1$ & $\lambda = 0.1$.

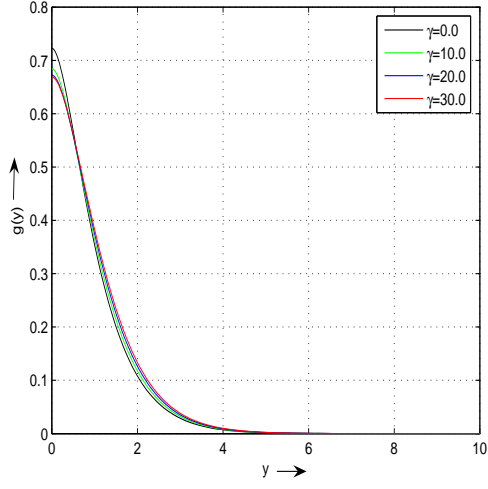


Fig. 18. Variation of g with γ at $\alpha_1 = 1$, $\beta = 2$, $V_0 = 0.5$, $M_n = 1$ & $\lambda = 0.1$.

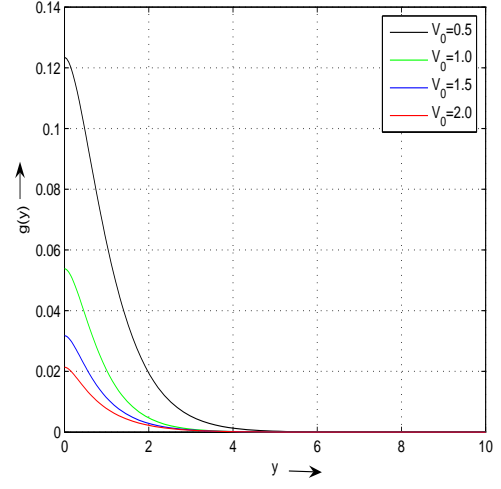


Fig. 19. Variation of g with V_0 at $\alpha_1 = 1$, $\beta = 2$, $\gamma = 1$, $M_n = 1$ & $\lambda = 1$.

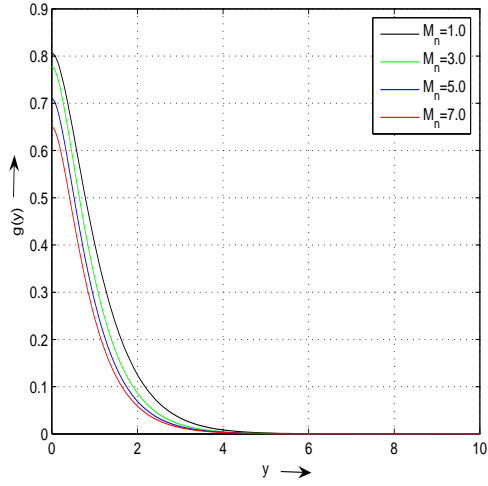


Fig. 20. Variation of g with M_n at $\alpha_1 = 1$, $\beta = 2$, $\gamma = 1$, $V_0 = 0.5$ & $\lambda = 0.5$.

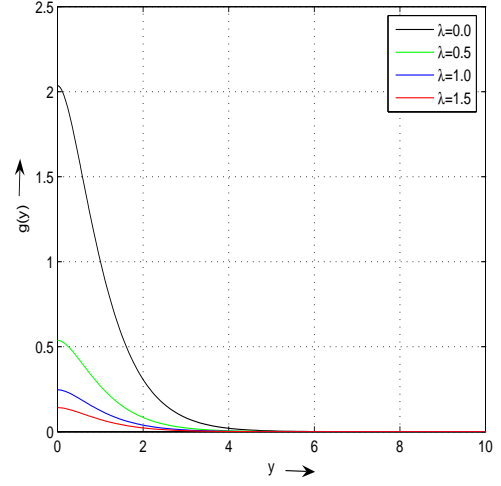


Fig. 21. Variation of g with λ at $\alpha_1 = 1$, $\beta = 2$, $\gamma = 1$, $V_0 = 0.5$ & $M_n = 1$.

α_1	β	γ	V_0	M_n	λ	P_r	E_c	$\theta'(0)$	$g(0)$
0.1								-3.826486	0.043378
0.3	1.0	2.0	2.0	1.0	1.0	2.0	1.0	-3.827067	0.043233
0.5								-3.827638	0.043090
	0.5							-2.840175	0.053275
1.0	1.5	2.0	1.0	1.0	1.0	3.0	0.5	-2.839138	0.053621
	2.5							-2.838103	0.053966
		1.0						-0.490230	0.026147
2.0	1.0	1.5	0.5	1.0	2.0	1.0	0.3	-0.490232	0.026142
		2.0						-0.490235	0.026137
			1.0					-0.492547	0.021544
3.0	2.0	1.0	3.0	0.5	1.0	3.0	0.5	-0.503001	0.000777
			5.0					-0.503356	0.000071
				0.1				-0.935282	0.064760
3.0	1.0	2.0	1.0	0.3	1.0	3.0	0.5	-0.969405	0.030639
				0.5				-0.969405	0.030639
					1.0			-2.848080	0.050640
4.0	2.0	2.0	1.0	1.0	2.0	3.0	0.5	-2.949012	0.016996
					3.0			-2.974725	0.008425
						1.0		-1.958481	0.020760
3.0	1.0	2.0	2.0	1.0	1.0	2.0	0.5	-3.916961	0.020759
						3.0		-5.875442	0.020757
							0.5	-3.973168	0.006708
1.0	3.0	2.0	2.0	1.0	2.0	2.0	1.5	-3.919504	0.020124
							2.0	-3.865840	0.033540

Table 1

Variations of $\theta'(0)$ and $g(0)$ with different flow parameters.